

# GENERALIZED GEOMETRY IN GRAVITY

## RTG “MODELS OF GRAVITY” WORKSHOP

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February 21, 2018

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## ALGEBROIDS

**Leibniz algebroid:**  $(E, \rho, [\cdot, \cdot])$ ,

$E \xrightarrow{\pi} M$ ,  $\rho \in \text{Hom}(E, TM)$  “anchor”,  $[\cdot, \cdot] : \Gamma(E) \times \Gamma(E) \rightarrow \Gamma(E)$

- 1  $[e, [e', e'']] = [[e, e'], e''] + [e', [e, e'']]$   $e, e', e'' \in \Gamma(E)$
- 2  $[e, fe'] = (\rho(e)f)e' + f[e, e']$ ,  $f \in C^\infty(M)$

**Lie algebroids:** Leibniz algebroids with skew-symmetric brackets

$$\longrightarrow \exists d_E : \Omega^\bullet(E) \rightarrow \Omega^{\bullet+1}(E)$$

**Courant algebroid:**  $(E, \rho, [\cdot, \cdot], \langle \cdot, \cdot \rangle)$ , Leibniz algebroid

$\langle \cdot, \cdot \rangle : \Gamma(E) \times \Gamma(E) \rightarrow C^\infty(M)$  “pairing” or  $g_E$ ,

- 1  $\rho(e)\langle e', e'' \rangle = \langle [e, e'], e'' \rangle + \langle e', [e, e''] \rangle \implies \mathcal{L}_e g_E = 0$
- 2  $\langle [e, e], e' \rangle = \frac{1}{2}\rho(e')\langle e, e \rangle$ .

## BRACKETS

$E$  exact Courant algebroid:  $0 \rightarrow T^*M \xrightarrow{j} E \xrightarrow{\rho} TM \rightarrow 0$ ,  
 $j := g_E^{-1} \circ \rho^T$ ,  $\text{Im} j = \ker \rho$ ,

(Ševera) unique (up to an isomorphism) classification via  $[H] \in H_3(M, \mathbb{R})$ :

$\exists$  isotropic splitting  $s: TM \rightarrow E$ ,  $\langle s(X), s(Y) \rangle = 0$  s.t.:

$$[s(X) + j(\xi), s(Y) + j(\eta)]_E = s([X, Y]_{\text{Lie}}) + j(\mathcal{L}_X \eta - i_Y d\xi - H(X, Y, \cdot))$$

$$\implies (E, \rho, [\cdot, \cdot]_E, \langle \cdot, \cdot \rangle_E) \cong (TM \oplus T^*M, pr_{TM}, [\cdot, \cdot]_D^H, \langle \cdot, \cdot \rangle_{TM \oplus T^*M}).$$

Twisting: if  $H \in \Omega_{\text{closed}}^3(M)$ ,  $B \in \Omega^2(M)$ ,  $\exp(B) = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \in O(d, d)$ :

$$\implies [\exp(B)e, \exp(B)e']_D^H = \exp(B) \left( [e, e']_D^{H+dB} \right)$$

$$i_X i_Y H(Z) \equiv \langle [s(X), s(Y)], s(Z) \rangle$$

- **Dorfman** bracket:

$$[V, W]_D := [X, Y]_{\text{Lie}} + \mathcal{L}_X \eta - i_Y d\xi.$$

- **Courant** bracket:  $[V, W]_C := \frac{1}{2} ([V, W]_E - [W, V]_E)$

$$V, W \in TM \oplus T^*M$$

$$V = X + \xi$$

$$W = Y + \eta$$

(I APPROACH:) GENERALIZED METRIC  $\mathbf{G}_\tau$ 

for  $\tau \in \text{End}(E)$ ,  $\tau^2 = 1$  (involution),

$$\mathbf{G}_\tau(V, W) := \langle V, \tau(W) \rangle$$

**Theorem:** if  $E$  vector bundle with fiber-wise metric  $\langle \cdot, \cdot \rangle_E$  of constant signature  $(p, q)$ ,  $\implies$  def. of  $\mathbf{G}_\tau$  is equivalent to def. of positive  $C_+ \subseteq E$  of rank  $p$

**Theorem:** if  $L, L^*$  isotropic subbundles of  $E$ ,  $E = L \oplus L^*$ , with  $\langle \cdot, \cdot \rangle$  of signature  $(n, n) \implies \mathbf{G}_\tau$  is equivalent to unique  $g \in \Gamma(S^2 L^*)$  and  $B \in \Omega^2(L)$

$$\mathbf{G}_\tau = \begin{pmatrix} g - Bg^{-1}B & Bg^{-1} \\ -g^{-1}B & g^{-1} \end{pmatrix}, \text{ for } L = TM, L^* = T^*M$$

$$\mathbf{G}_\tau \in \text{O}(d, d) / (\text{O}(d) \times \text{O}(d))$$

for a more general reduction:

$$\text{O}(t, s) \supset \text{O}(p, q) \times \text{O}(t-p, s-q)$$

 $\implies$ 

- $C_+ \cap T^*M = 0$

- $\text{rk} C_+ = \text{rk} E - \dim M$

# (I APPROACH:) BISMUT CONNECTION $\tilde{\nabla}^0$

in extended tangent space:

- $\nabla$  Courant connection if  $\nabla_e(fe') = f\nabla_e e' + \rho(e)(f)e'$  and  $\langle \nabla_e e', e'' \rangle + \langle e', \nabla_e e'' \rangle = \rho(e)(\langle e', e'' \rangle)$
- torsion  $T_\nabla \in \mathcal{T}_3$  e.g.  $T_\nabla(e, e', e'') := \langle \nabla_e e' - \nabla_{e'} e - [e, e'], e'' \rangle + \langle \nabla_{e''} e, e' \rangle$
- extended Riemann tensor, if  $R_\nabla^{(0)}(k', k, e, e') := \langle R(e, e')k, k' \rangle$ ,  
 $R_\nabla(k', k, e, e') := \frac{1}{2} \left\{ R_\nabla^{(0)}(k', k, e, e') + R_\nabla^{(0)}(e', e, k, k') + \langle \nabla_{e_\lambda} e, e' \rangle_E \cdot \langle \nabla_{k_\lambda} k, k' \rangle \right\}$

one amongst the possible generalizations of the std Levi-Civita connection (arXiv:1512.08522, Jurco, Vysoky):

$$\tilde{\nabla}_X^0 = \begin{pmatrix} \nabla_X^{LC} + \frac{1}{6}g^{-1}H'(g^{-1}(\xi), \star, \cdot) & -\frac{1}{3}g^{-1}H'(X, g^{-1}(\star), \cdot) \\ -\frac{1}{3}H'(X, \star, \cdot) & \nabla_X^{LC} + \frac{1}{6}H'(g^{-1}(\xi), g^{-1}(\star), \cdot) \end{pmatrix}.$$

$$\text{Ric}_{\tilde{\nabla}^0}(e, e') := R_{\tilde{\nabla}^0}(e^\lambda, e, e_\lambda, e'), \quad R_{\tilde{\nabla}^0}^{\mathbf{G}} := \text{Ric}_{\tilde{\nabla}^0}(\mathbf{G}^{-1}(e^\lambda), e_\lambda),$$

## (II APPROACH:) RECAP OF GRADED GEOMETRY

( $\mathbb{Z}$ -)graded vector space  $A$ :  $A = \bigoplus_{i \in \mathbb{Z}} A_i$ , over a field  $k$  of characteristic zero  
 $\iff$  graded Poisson algebra of degree  $n$ ,  $\{\cdot, \cdot\}$  of degree  $-n$ ,  $\cdot$  zero-graded commutative product:

- $\{a, b\} = -(-1)^{(|a|+n)(|b|+n)} \{b, a\}$ ,
- $\{a, \{b, c\}\} = \{\{a, b\}, c\} + (-1)^{(|a|+n)(|b|+n)} \{a, \{b, c\}\}$   $|\cdot| := \text{degree}$
- $\{a, b \cdot c\} = \{a, b\} \cdot c + (-1)^{|b|(|a|+n)} b \cdot \{a, c\}$

Graded manifold  $M$ : its local coordinates have a degree  $|\cdot|$

Graded vector bundle  $E \xrightarrow{\pi} M$ : coords  $\{x\}$  of deg 0 for  $M$ , coords  $\{v\}$  of deg  $n$  for  $E \implies$  its algebra of sections respects the grading

“coords even (in parity) have even degree, odd ones have odd degree”  $\iff$   
*“N-manifold”*

when  $\exists Q \in \Gamma(TM)$ , of degree 1, s.t.  $[Q, Q] = 2Q^2 = 0$ ,  $\implies$  *“NQ-manifold”*

to this setup, add a symplectic structure, i.e. a closed non-degenerate 2-form  $\Omega$  of degree  $n$ , “Hamiltonian”



## (II APPROACH:) POISSON GRADED MANIFOLDS &amp; COURANT ALGEBROIDS

**Theorem:** *Symplectic  $N$ -manifolds of degree 2 are 1:1 with pseudo-Euclidean vector bundles* [Roytenberg, arXiv:math/0203110v1]

$(M, \Omega)$  with  $n = 2$  has homological vector field  $Q$  given by  $\Theta$  of degree 3,  
 $\{\Theta, \Theta\} = 0$ ,

$$Q := \{\Theta, \cdot\}$$

## COURANT ALGEBROIDS SUPPORT GRADING

**Theorem:** symplectic  $NQ$ -manifolds of degree 2 are 1 : 1 with Courant algebroids,

in particular,  $(T^*[2]T[1]M, pr_{TM}, [\cdot, \cdot]_D, \langle \cdot, \cdot \rangle)$  is a Courant algebroid

**Derived brackets :**

$$\{\{\cdot, \Theta\}, \cdot\}$$

$$\{\{e, \Theta\}, f\} =: \rho(e) \cdot f$$

$$\{\{e, \Theta\}, e'\} =: [e, e']_{\text{Dorfman}}$$

$\{\cdot, \cdot\}$  is of degree -2, hence even

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$(g_{\mu\nu}, B_{\mu\nu}, \phi)$  FIELDS IN STRING THEORY (WITH APPROACH II)

base manifold  $M$ ,  $T^*[2]T[1]M$  total space,  $\{x, \xi, \theta, p\}$

$C^\infty(M) \ni f = f(x)$ ,

$\Gamma(T^*[2]M) \ni v, \omega = \omega^j p_j$        $\Gamma(T[1]M \oplus T^*[1]M) \ni V, W = Y^i \chi_i + \eta_i \theta^i \rightarrow \xi_\alpha := (\chi_i, \theta^i), W = W^\alpha \xi_\alpha, \tilde{\xi}^\beta$  dual

Poisson graded structure given by:

canonical Poisson brackets

deformation

$$\left\{ \begin{array}{l} \{g, f\} = 0 \\ \{p_i, f\} = \partial_i f \\ \{p_i, \xi_\alpha\} = 0 \\ \{\xi_\alpha, \xi_\beta\} = \eta_{\alpha\beta} \\ \{p_i, p_j\} = 0 \end{array} \right. \implies \left\{ \begin{array}{l} \{g, f\} = 0 \\ \{p_i, f\} = \partial_i f \\ \{p_i, \xi_\alpha\} = \nabla_i \xi_\alpha \\ \{\xi_\alpha, \xi_\beta\} = G_{\alpha\beta}(x) \\ \{p_i, p_j\} = R_{ij} = 0 \end{array} \right.$$

$$G = \begin{pmatrix} 2g(x) & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ g(x) - B(x) & 1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ g(x) + B(x) & 1 \end{pmatrix}$$

$G$  corresponds to  $e^{(g+B)} \chi_k = \chi_k + (g_{kl} + B_{kl})\theta^l$

it is possible to canonically associate a cubic hamiltonian,  $\Theta$ , through the homological vector field as  $Q = \{\Theta, -\}$ ,

$$\Theta = \tilde{\xi}^\alpha \rho(\xi_\alpha) + \frac{1}{3!} C^{\alpha\beta\gamma}(x) \xi_\alpha \xi_\beta \xi_\gamma,$$

for  $\rho : E \rightarrow TM$  anchor,  $C^{\alpha\beta\gamma}$  3-tensor ("fluxes").

## (INTERMEDIATE) RESULTS

given that  $\nabla_{\rho(\xi_\alpha)} \xi_\beta = \mathbb{W}_{\alpha\beta}{}^\gamma \xi_\gamma = \{\rho(\xi_\alpha), \xi_\beta\}$ ,  $\mathbb{W} = \partial(g+B)$

$$\begin{aligned} \{\{V, \Theta\}, W, U\} &= \langle \nabla_V W, U \rangle - \langle \nabla_W V, U \rangle + \langle \nabla_U V, W \rangle + C(V, W, U) \\ &=: \langle [V, W]', U \rangle \end{aligned}$$

$$\begin{aligned} [V, W]' &= [\rho(V), \rho(W)]_{\text{Lie}} + T(V, W) + \langle \nabla \cdot V, W \rangle + \text{“fluxes”} \\ &=: [\rho(V), \rho(W)]_{\text{Lie}} + \langle \tilde{\nabla} V, W \rangle + \text{“fluxes”} \end{aligned}$$

$\langle [V, W]', U \rangle$  is Koszul formula for  $g+B$ :

$$\begin{aligned} 2g(\tilde{\nabla}_X Y, Z) &= g([X, Y], Z) - g([X, Z], Y) - \\ g([Y, Z], X) &+ \partial_X(g(Y, Z) + B(Y, Z)) + \\ \partial_Y(g(X, Z) + B(X, Z)) &- \partial_Z(g(X, Y) + B(X, Y)) \end{aligned} \iff \begin{aligned} \mathbb{W}_{\alpha\beta\gamma} - \mathbb{W}_{\beta\alpha\gamma} + \mathbb{W}_{\gamma\alpha\beta} &= \\ 2\Gamma_{\beta(\gamma\alpha)} + H_{\alpha\beta\gamma} \end{aligned}$$

↓

gravity sector of string theory for dilaton  $\phi = 0$  reconstructed from Ricci tensor  $R_{\mu\nu}$  for generalized  $T^*[1]M \oplus T[1]M$ , contracted with  $(g^{\mu\nu} + B^{\mu\nu})$ :

$$\mathcal{S}_{\sigma\text{-st.thy}} = \frac{1}{4\pi\kappa^2} \int \text{Vol}(g) \left( R - \frac{1}{12} H^2 \right)$$

## (INTERMEDIATE) RESULTS

To obtain the  $\sigma$ -model for the full closed string sector (26-dim for bosonic string, 10-dim for RNS-string), seek for deformations whose connection is given by

$$\Gamma_{\nu\delta}^{\text{L.C.}\beta} + \frac{1}{2}H_{\nu\delta}^{\beta} - \frac{2}{d-1}\partial_{\delta}\phi\delta_{\nu}^{\beta},$$

reconstruction of  $\mathcal{L}_{\sigma\text{-st.thy}}$  goes as before

$$\begin{aligned} \implies (g^{\mu\nu} + B^{\mu\nu})e^{-2\phi} \left( \beta_{\mu\nu}^g + \beta_{\mu\nu}^B - \frac{4}{d-1}\partial_{\mu}\phi\partial_{\nu}\phi \right) \\ \mathcal{L} = e^{-2\phi} \left( R - \frac{1}{12}H^2 + 4(\partial\phi)^2 \right) \end{aligned}$$

where the beta functions are:

$$\beta_{\mu\nu}^g + \beta_{\mu\nu}^B = R_{\mu\nu}^{\text{L.C.}} - \frac{1}{4}H_{\mu\lambda\varepsilon}H_{\nu}^{\lambda\varepsilon} + 2\nabla_{\mu}\partial_{\nu}\phi - \frac{1}{2}\nabla^{\lambda}H_{\lambda\mu\nu} + \partial^{\lambda}\phi H_{\lambda\mu\nu}$$

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### 3. FERMIONIC T-DUALITY

We argue that transformations (fermionic T-dualities) of string fluxes could be related to a deformation of a Courant algebroid and its bracket.

$$\mathcal{L}_\sigma^{GS} = \int d^2z B_{00}(Y) \partial\theta^0 \bar{\partial}\theta^0 + L_{0M}(Y) \partial\theta^0 \bar{\partial}Y^M + L_{M0}(Y) \partial Y^M \bar{\partial}\theta^0 + L_{MN}(Y) \partial Y^M \bar{\partial}Y^N,$$

isometry  $\theta^0 \rightarrow \theta^0 + c \implies$  conservation in one direction ( $\exists$  a killing spinor). Integrating it out (by means of its e.o.m.'s) one ends with the same type of theory IIA (or IIB) if the  $F^i$  R-R fields and the dilaton  $\phi$  change according to

$$\exp(\phi') F' = \exp(\phi) F \pm 32 \sum_{I,J=1}^N (\varepsilon_I \otimes \hat{\varepsilon}_J) M_{IJ}$$

$$\phi' = \phi + \frac{1}{2} \text{Tr}(\log M)^{-1}$$

$M^{-1}$  given by:

$$\text{and } \begin{pmatrix} \varepsilon \\ \hat{\varepsilon} \end{pmatrix}$$

$$\partial_a (M^{-1})_{IJ} = i\varepsilon_I \gamma_a \varepsilon_J \mp i\hat{\varepsilon}_I \gamma_a \hat{\varepsilon}_J,$$

SO(2)-invariant Majorana-Weyl spinor (same chirality)

## 4. κ-SYMMETRY

We argue that κ-symmetry of RNS string theory could have a geometrical meaning in the generalized geometry setup.

Spacetime supersymmetric Nambu-Goto string action:

$$\mathcal{S}_{NG1} = -\frac{\alpha'}{\pi} \int d^2\sigma \sqrt{-\det [(\partial_\alpha X^\mu - \bar{\theta}^A \Gamma^\mu \partial_\alpha \theta^A) (\partial_\beta X_\mu - \bar{\theta}^A \Gamma_\mu \partial_\beta \theta^A)]},$$

global supersymmetry:  $\delta \theta^{Aa} = \varepsilon^{Aa}$ ,  $\delta X^\mu = \bar{\varepsilon}^a \Gamma^\mu \theta^A$ , with algebra  $[\delta_1, \delta_2] \theta = 0$ ,  $[\delta_1, \delta_2] X^\mu = -2\bar{\varepsilon}_1 \Gamma^\mu \varepsilon_2$ .

Equations of motion:  $\dot{P} = 0$ ,  $\dot{\theta} = 0 \implies$  only half of the dofs of  $\theta$  are genuine.

$$\implies \mathcal{S}_{WZ2} = -m \int d\tau \bar{\theta} \Gamma_{11} \partial_\tau \theta$$

is needed.

This preserves a hidden local symmetry called κ-symmetry.



# SUMMARY

- *generalized geometry: Courant algebroids for a “doubled” total space  $TM \oplus T^*M$ , usually extended Riemannian geometry is assigned;*
- *1:1 correspondance with graded symplectic manifolds of deg 2, endowed with  $Q$  homological vector field, and  $\Theta$  hamiltonian*
  - aim: a full derivation of closed strings in low-energy string theory  $\leftarrow$  so far, we have a connection, still a good deformation of the Poisson bracket, s.t.  $\phi$  is included, is missed;
  - aim: a study of the transformation of fluxes in supersymmetric string theory: fermionic T-duality &  $\kappa$ -symmetry

# Thanks for your attention