Quantum gravity in the asymptotic-safety approach

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RTG Colloquium Jacobs University Bremen

22 April 2020



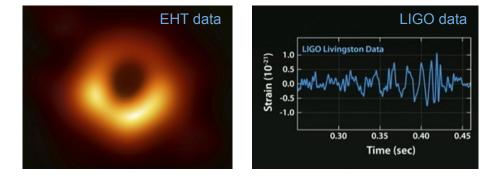
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Einstein Gravity, Singularities and Renormalization

Einstein Gravity

$$S_{EH} = rac{1}{16\pi G}\int d^4x \sqrt{-g}\;(R-2\Lambda)$$



Very successful in the description of gravitational phenomena, even in strong field regimes. However:

- Classical singularities

the solutions of the Einstein equations are generally singular \Rightarrow loss of predictivity

[Penrose, '65; Penrose, Hawking, '70;]

- Renormalizability issues

the Newton's constant has negative mass dimension \Rightarrow ultraviolet divergences

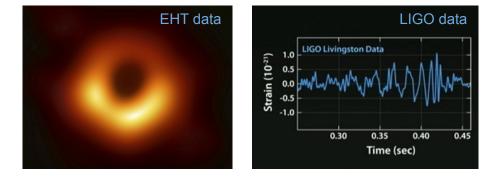
['t Hooft, Veltman, '74; Goroff, Sagnotti, '86;]

General Relativity is incomplete and breaks down at high energies

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Asymptotically Safe Gravity (Weinberg, 1976)

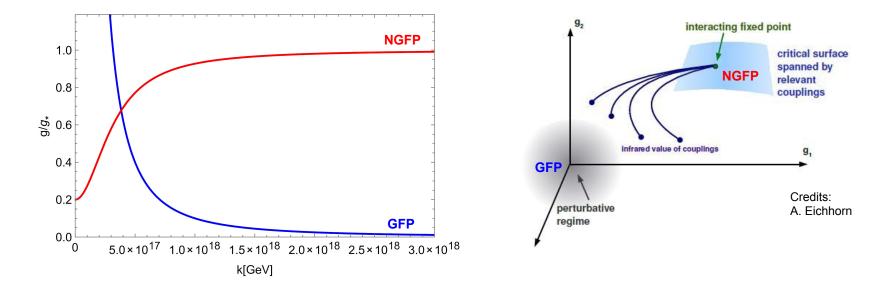
- General Relativity might make sense as a quantum field theory
- Key idea: generalized notion of renormalizability based on the Wilsonian Renormalization Group

S. Weinberg, Erice Subnucl. Phys.1976:1

RG fixed points and renormalizability

Two types of well-definite ultraviolet completion (microscopic/fundamental theory)

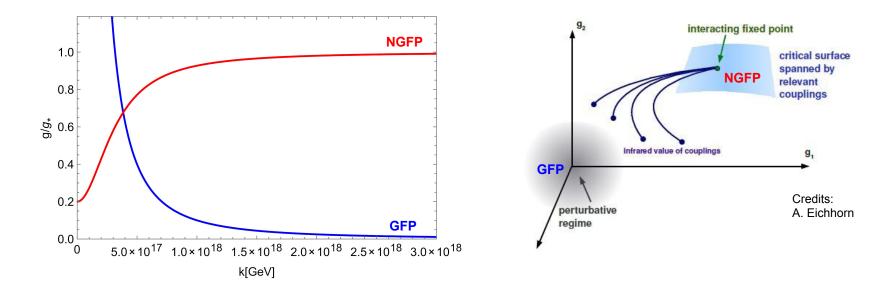
- Gaussian Fixed Point (GFP): free theory ⇒ Asymptotic Freedom
- Non-gaussian Fixed Point (NGFP): interacting theory ⇒ Asymptotic Safety



RG fixed points and renormalizability

Generalized (non-perturbative) renormalizability

- Ultraviolet completion ⇔ UV-attractive fixed point (microscopic theory)
 - Gaussian Fixed Point (GFP): free theory ⇒ Asymptotic Freedom
 - Non-gaussian Fixed Point (NGFP): interacting theory ⇒ Asymptotic Safety
- **Predictivity** ⇔ finite number of relevant directions



Asymptotic Safety in Quantum Gravity

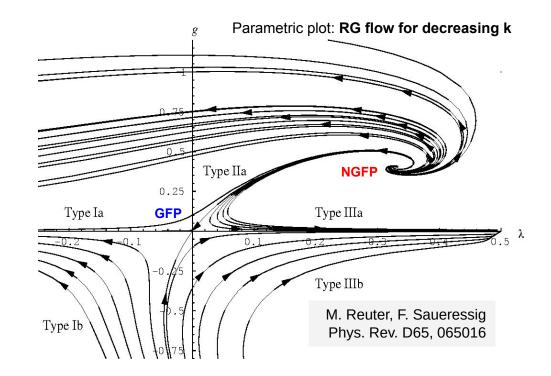
Einstein-Hilbert truncation

$$S_k = rac{1}{16\pi G_k}\int d^4x \sqrt{-g}\;(R-2\Lambda_k)$$

$$G_k = k^{-2} g_k \qquad \Lambda_k = k^2 \lambda_k$$

Fixed points of the RG flow:

- **GFP** \rightarrow saddle point
- **NGFP** \rightarrow UV-attractor



Looking for Asymptotic Safety

Functional RG equations for different ansatz of S allow to verify the existence of the NGFP

$$k\partial_k\Gamma_k=rac{1}{2}\mathrm{STr}\left\{\left(\Gamma_k^{(2)}+\mathcal{R}_k
ight)^{-1}\;k\partial_k\mathcal{R}_k
ight\}$$

C. Wetterich. *Phys. Lett. B* 301:90 (1993) M. Reuter. *Phys. Rev.* D. **57** (2): 971 (1998)

→ NGFP + 3 relevant directions

$f(R)\simeq \sum_{n=0}a_nR^n$.	$f(R) \simeq$	
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Polynomial up to N=71:

Reuter, Lauscher, '02; Codello, Percacci, Rahmede '09; Benedetti, Caravelli, '12; Dietz, Morris, '12; Falls, Litim, Nikolakopoulos, Rahmede, '13, '14 Demmel, Saueressig, Zanusso, '15; Falls, Litim, Schoeder, '18 **Beyond polynomial**: Benedetti, Caravelli, '12; Demmel, Saueressig, Zanusso, '12; Dietz, Morris, '13

$$R,~R^2,~R_{\mu
u}R^{\mu
u}$$

Benedetti, Machado, Saueressig, '09. Christiansen, '16, Oda, Yamada '17

 $R, \ C^{\kappa\lambda}_{\mu\nu}C^{\rho\sigma}_{\kappa\lambda}C^{\mu\nu}_{\rho\sigma}$

Gies, Knorr, Lippoldt, Saueressig '16

→ Canonical power counting is still a good guideline Eichhorn, Lippoldt, Pawlowski, Reichert, Schiffer. Phys.Lett. B792 (2019) 310-314

Asymptotically Safe Gravity: problems and open questions

Main problems related with the Functional RG:

$$k\partial_k\Gamma_k=rac{1}{2}{
m STr}\left\{\left(\Gamma_k^{(2)}+{\cal R}_k
ight)^{-1}\,k\partial_k{\cal R}_k
ight\}$$

• Truncation

- Solving the FRG equations requires truncation
- Within truncation: regulator, gauge and parametrization dependence

Form factors, RG flow of generic entire functions: Knorr, Ripken, Saueressig - Class. Quant. Grav. 36 (2019) 23, 234001

• Causal structure and Lorentzian signature

Lorentzian EH truncation, Matsubara formalism: Manrique, Rechenberger, Saueressig - *Phys.Rev.Lett.* 106 (2011) 25130 Euclidean EH truncation, ADM formalism: Biemans, <u>AP</u>, Saueressig - *Phys.Rev.D* 95 (2017) 8, 086013 Euclidean EH truncation + matter, ADM formalism: Biemans, <u>AP</u>, Saueressig - *JHEP* 05 (2017) 093 Euclidean EH truncation, generic foliation, lorentz-violation: Knorr *Phys.Lett.B* 792 (2019) 142-148 Euclidean EH truncation + matter, generic foliation, lorentz-violation: Eichhorn, <u>AP</u>, Schiffer. Arxiv:1911.10066

More open questions on asymptotically safe gravity:

• Unitarity

Eichhorn, <u>AP</u>, Wetterich - to appear soon

Phenomenology / signatures of asymptotic safety

Literature on RG-improved models: A. Bonanno, A. Eichhorn, K. Falls, A. Held, D. Litim, AP, M. Reuter, F. Saueressig etc.

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Asymptotic Safety on foliated spacetimes

- Functional Renormalization Group is based on a *Euclidean* path integral
- The nature of our universe is Lorentzian
- QFTs on fixed Minkowski background: Wick rotation to restore a *Lorentzian* signature
- **Gravity**: integration over all possible geometries

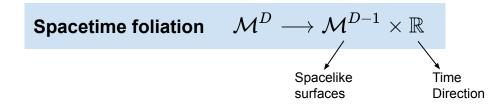
$$\int D[g_{\mu
u}]\,e^{-S[g_{\mu
u}]}$$
 \longrightarrow

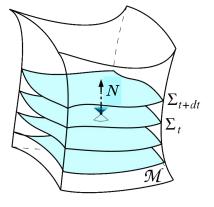
Metric formulation: no well-distinguished time direction

- The **causal structure** of the spacetime might be important to determine the quantum properties of gravity
- Comparison of FRG results with causal dynamical triangulations (CDT) program

R. Loll. Nucl.Phys.Proc.Suppl. 94 (2001) 96-107

Asymptotic Safety on foliated spacetimes



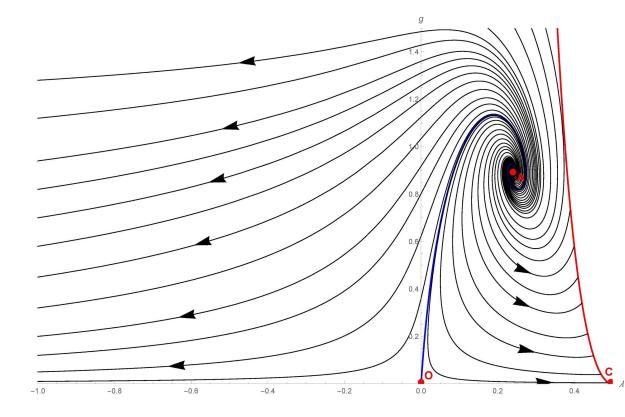


The covariant metric is written in terms of the ADM fields

Functional integral over the ADM-fields

$$\mathcal{Z} = \int D[N] \, D[N^i] \, D[\sigma_{ij}] \, e^{-S[N,N_i,\sigma_{ij}]} \quad \Longrightarrow \quad \Gamma_k = \Gamma_k[N,N_i,\sigma_{ij}]$$

Properties of the gravitational RG flow: phase diagram in (3+1)-dimensions



The RG-flows obtained in the covariant and foliated approaches share similar properties:

- UV-attractive NGFP
- Complex critical exponents

J. Biemans, <u>AP</u>, F. Saueressig Phys. Rev. D 95, 086013 (2017)

RG-flow: Gravity-matter systems

$$R_{\mu
u} - rac{1}{2} R \, g_{\mu
u} = 8 \pi G \, T_{\mu
u}$$

RG-flow: Gravity-matter systems

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u} & = 8\pi G\,T_{\mu
u} \ & \downarrow & \downarrow \ & \downarrow & \downarrow \ & S_k & = S_k^{ ext{grav}} + rac{S_k^{ ext{matter}}}{k} \end{aligned}$$

Is the gravitational RG flow influenced by the

presence of matter fields?

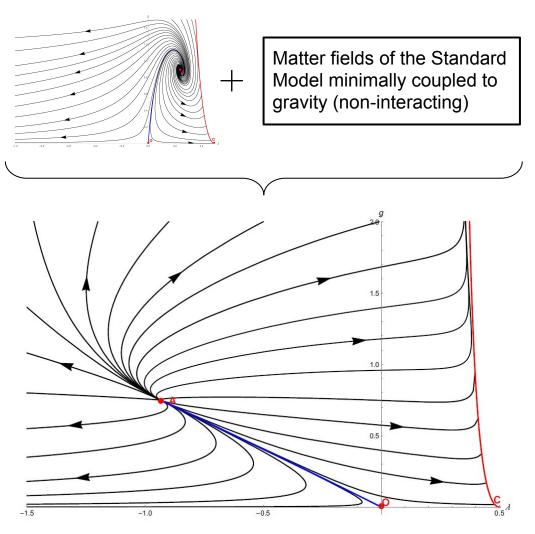
Matter...matters!

 Adding a minimal coupling to matter fields modifies *universal properties* of the gravitational RG flow

> P. Dona, A. Eichhorn, R. Percacci Phys. Rev. D **89**, 084035 (2014)

 Vice versa, gravity influences the running of the matter couplings and could induce a viable UV completion of the SM

> Daum, Harst, Reuter '09; Harst, Reuter '11; Folkerts, Litim, Pawlowski '11; Christiansen, Eichhorn '17; Eichhorn, Held, Pawlowski '16; Eichhorn, Held '17



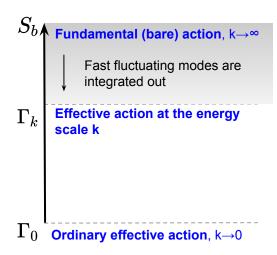
J. Biemans, <u>AP</u>, F. Saueressig JHEP 05 (2017) 093 (2017)

Astrophysical and cosmological implications of Asymptotically Safe Gravity

Goal: understand how *quantum fluctuations* modify the classical solutions of GR at high energies

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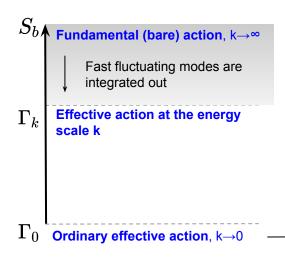
Derivation of quantum solutions from the functional renormalization group:



$$k\partial_k\Gamma_k=rac{1}{2}\mathrm{STr}\left\{\left(\Gamma_k^{(2)}+\mathcal{R}_k
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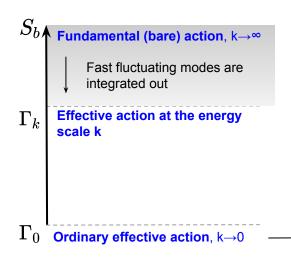
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All quantum fluctuations are integrated out (non-localities) Incorporates all quantum effects \rightarrow fully-dressed quantities Quantum solutions from effective field equations:

$$rac{\delta\Gamma_0}{\delta g_{\mu
u}}=0$$

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$${\delta \Gamma_0 \over \delta g_{\mu
u}} = 0$$

Problem: Fully-quantum effective action is currently unknown

Phenomenology of AS:

- Fully-quantum effective action is unknown
- Semi-classical models that capture key features / hallmarks of asymptotic safety

Hallmarks of asymptotically safe gravity:

- **Scale invariance** of the theory at high energies
- **Gravitational antiscreening** (weakening of the gravitational interaction at high energies / short distances)

Idea:

- \Rightarrow Construct *semi-classical models* that have these features
- ⇒ In the context of quantum field theory: renormalization-group improvement (expected to provide qualitative understanding of the effect of quantum fluctuations)

The Renormalization Group improvement

- Asymptotically Safe Gravity is based on Quantum Field Theory
- In the context of QFT, the **RG-improvement** is widely used to study the effect of leading order quantum corrections
 - Start from a well-known classical system
 - Replace characteristic constants with running couplings
 - Close the system by identifying the RG-scale with a characteristic energy scale of the system

EXAMPLE 1: RG-improvement of the electric potential in *Quantum Electrodynamics* $V(r) = -\frac{e^2}{r} \rightarrow V_k(r) = -\frac{e_k^2}{r} \xrightarrow[k \sim 1/r]{} V_1(r) \sim -\frac{e^2(r_0^{-1})}{4\pi r}(1 + b\log(r_0/r))$ 1-loop Uehling potential

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EXAMPLE 2: RG-improvement of the action in Scalar Electrodynamics Radiative corrections included by identifying the RG-scale with the field strength $S_{cl}[\phi] \longrightarrow S_k[\phi] \quad k \sim \phi \quad \Rightarrow \quad \text{Coleman-Weinberg effective potential}$

⇒ Phenomenology of AS:

Semi-classical models that capture key features / hallmarks of asymptotic safety Models in the literature based on RG-improvement procedure

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\Rightarrow Phenomenology of AS:

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Hallmarks of asymptotically safe gravity:

- **Scale invariance** of the theory at high energies

Explanation for the nearly-scale invariant power spectrum of temperature fluctuations in the CMB?

Gravitational antiscreening (weakening of the gravitational interaction at high energies / short distances)

Weakening / resolution of classical spacetime singularities?

Modeling the gravitational antiscreening with the RG-improvement

Renormalization group equations ⇒ **running Newton's coupling** •

 $G_k = rac{G_0}{1+G_0 \; q_\star^{-1} k^2}$ A. Bonanno, M. Reuter Phys. Rev. D 62, 043008 (2000)

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If the Asymptotic Safety conjecture holds, there exists a scale-invariant regime at high energies • where the Newton coupling scales as

 $G_k = g_* k^{-2}$ for $k \gg M_{Pl}$

 \Rightarrow the RG-scale-dependent Newton coupling vanishes at high energies

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The RG-scale-dependence of the Newton's constant can be used to model the weakening of • the gravitational interaction at high energies

⇒ weakening of the singularities typically appearing in the classical theory

A. Bonanno, M. Reuter. PRD 62 (2000) 043008 A. Bonanno, B. Koch, AP CQG 34 (2017) 095012 A. Bonanno, G. Gionti, AP, CQG 35 (2018) 6, 065004 AP. EPJC 79 (2019) 470

- Let us assume that **inflation** occurs in the presence of gravity (EH approximation), minimally coupled to matter fields
- In a first approximation, we assume that matter does not contribute to the inflationary dynamics, i.e. inflation is driven by QG fluctuations only
- 1) Running couplings induce a scale-dependent Lagrangian

$$\mathcal{L}_k = rac{1}{16\pi G_k} \{R-2\Lambda_k\} - eta_k R^2 + \dots$$

2) In the fixed-point regime diffeomorphism-invariance (contracted Bianchi identities) implies

$$Rigg(rac{1}{G_k}igg)' - 2igg(rac{\Lambda_k}{G_k}igg)' = 0 \qquad \Rightarrow \qquad k^2 \sim R$$

M. Reuter, H. Weyer, JCAP 0412, 001 (2004)
B. Koch, I. Ramirez, CQG 28 (2011) 055008
S. Domazet, H. Stefancic, CQG 29 (2012) 235005

 \Rightarrow Effective f(R) action

- The critical scaling around the NGFP depends on the matter content of the theory
- Gravity+matter → positive and real critical exponents

J. Biemans, <u>AP</u>, F. Saueressig JHEP 05 (2017) 093 (2017)

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Einstein-Hilbert truncation

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The RG-improved Lagrangian at the fixed point is an effective f(R) theory:

 $f_*(R) \sim R^2$

Consistent with quadratic fixed-point solution from the RG flow of f(R)-gravity

J. A. Dietz, T. R. Morris - JHEP01(2013)108 M. Demmel, F. Saueressig, O. Zanusso - JHEP08(2015)113

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$$egin{aligned} \mathcal{L}_k &= rac{k^2}{16\pi g_k} ig\{ R - 2\lambda_k k^2 ig\} & k^2 \sim R \ & ig| & igslessing ext{ Scaling in the vicinity} \ & igslessing ext{ of the NGFP} \ & igslessing \lambda_k &= \lambda_* + c_1 \, e_1^1 \, iggl(rac{k}{k_0} iggr)^{- heta_1} + c_2 \, e_2^1 \, iggl(rac{k}{k_0} iggr)^{- heta_2} \ & g_k &= g_* + c_1 \, e_1^2 \, iggl(rac{k}{k_0} iggr)^{- heta_1} + c_2 \, e_2^2 \, iggl(rac{k}{k_0} iggr)^{- heta_2} \ & iggr) \end{aligned}$$

⇒ The effective Lagrangian depends on universal properties on the RG flow

- Gravity+matter → positive and real critical exponents
- The effective Lagrangian depends on the critical exponents

$$\mathcal{L}_{ ext{eff}} = a_0 \, R^2 + b_1 \, R^{rac{4- heta_1- heta_2}{2}} + b_2 \, R^{rac{4- heta_1}{2}} + b_3 \, R^{rac{4- heta_2}{2}} + b_4 \, R^{2- heta_1} + b_5 \, R^{2- heta_2}$$

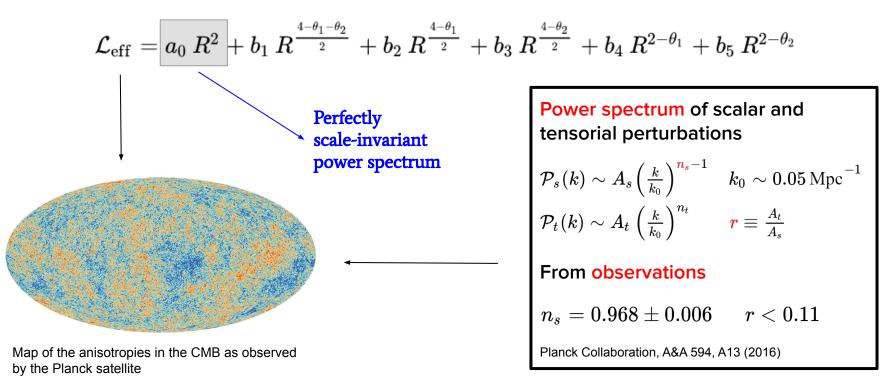
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$$\mathcal{L}_{eff} = a_0 R^2 + b_1 R^{\frac{4-\theta_1-\theta_2}{2}} + b_2 R^{\frac{4-\theta_1}{2}} + b_3 R^{\frac{4-\theta_2}{2}} + b_4 R^{2-\theta_1} + b_5 R^{2-\theta_2}$$

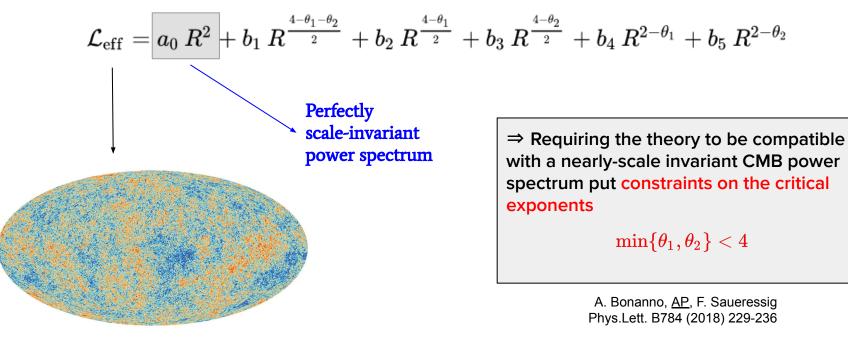
$$Power spectrum of scalar and tensorial perturbations
$$\mathcal{P}_s(k) \sim A_s \left(\frac{k}{k_0}\right)^{n_s-1} \quad k_0 \sim 0.05 \text{ Mpc}^{-1}$$

$$\mathcal{P}_t(k) \sim A_t \left(\frac{k}{k_0}\right)^{n_t} \quad r \equiv \frac{A_t}{A_s}$$
From observations
$$n_s = 0.968 \pm 0.006 \quad r < 0.11$$
Planck Collaboration, A&A 594, A13 (2016)$$

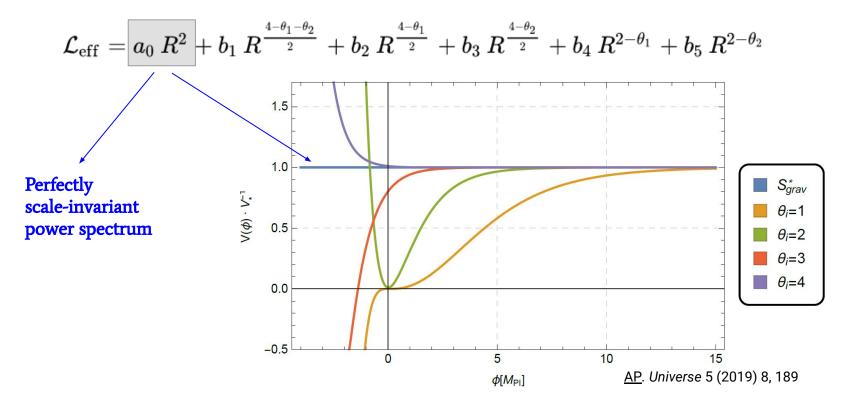
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Asymptotically Safe Inflation: the role of matter fields

A. Bonanno, <u>AP</u>, F. Saueressig Phys.Lett. B784 (2018) 229-236

Effective action from AS-gravity coupled to <u>SM-like matter fields</u>

$$\mathcal{L}_{ ext{eff}} = rac{1}{16\pi G_N} \Big(R + rac{1}{6m^2} R^2 - 2 \Lambda_{ ext{eff}} \Big)$$

Effective couplings

$$G = rac{g_*^2}{c_2 \left(2g_*\xi + 2\lambda_*\xi - 1
ight)M_{
m Pl}^2} \qquad m^2 = rac{g_*}{6 \, \xi \left(1 - 2\lambda_*\xi
ight)G} \qquad \Lambda_{
m eff} = rac{M_{
m Pl}^4 (c_2^2 + c_1 g_*)G}{g_*^2}$$

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⇒ The existence of an inflaton field requires the cosmological constant to be negative at the NGFP

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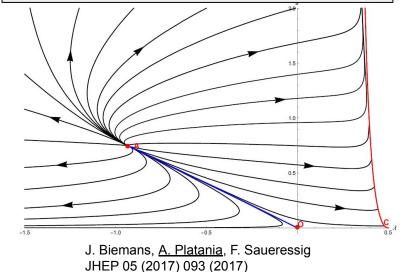
$$G = rac{g_*^2}{c_2 \ (2g_*\xi + 2\lambda_*\xi - 1) \ M_{
m Pl}^2} \qquad m^2 = rac{g_*}{6 \ \xi \ (1 - 2\lambda_*\xi) \ G}$$

⇒ The existence of an inflaton field requires the cosmological constant to be negative at the NGFP

This condition is compatible with similar constraints obtained in B. Koch and F. Saueressig, Class. Quantum Grav. 31, 015006 (2014) A. Eichhorn, A. Held. Phys. Lett. B777 (2018), pp. 217–221

A. Bonanno, G. Gionti, AP. Class. Quantum Grav. 35 (2018) 065004

Phase diagram: Einstein-Hilbert truncation + minimally-coupled matter fields of the Standard Model



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Summary

- <u>AS-gravity</u>: aims at constructing quantum gravity in a QFT framework Key ingredient: <u>NGFP</u> providing a well defined UV-completion for the gravitational interaction
- <u>Computations with matter</u>: matter fields influence the gravitational RG flow, and vice versa.
- <u>Computations on foliated spacetimes</u>: in the EH truncation qualitatively same phase diagram as in the standard case *Open questions*: Wick rotation / Lorentzian-signature computation?
- Phenomenology of asymptotically safe gravity: **RG-improved models of spacetimes**
- Hallmarks of asymptotically safe gravity:
 - Scale invariance of the theory at high energies (\Rightarrow nearly-scale invariant power spectrum?)
 - **Gravitational antiscreening** (\Rightarrow resolution of classical spacetime singularities?)
- More open questions and problems to be addressed
 - Truncation issue
 - Unitarity